

# Determination of the <sup>232</sup>Th nucleus number using small solid angle method

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### **1.1 Introducion**

 $^{232}$ Th: the most stable isotope of thorium (natural abundance 99.98%), is a fissionable nucleus which plays an important role in the Th/U fuel cycle.

 $^{232}$ Th (*n*, *f*) reaction: a important reaction in nuclear engineering applications and it could enhance our understanding in nuclear physics (such as "thorium anomaly" effect, the triple-humped barriers, etc).

### **1.2 Problems**

- I. Comparing to other fissile nuclei, such as  ${}^{238}U$  or  ${}^{235}U$ , the experimental measurements of  ${}^{232}Th$  are fewer for the (*n*, *f*) reaction.
- II. There is a systematic difference between the measurement results using different neutron sources (mono-energetic neutron sources and white neutron sources).

#### 1.3 Experimental data of <sup>232</sup>Th(*n*, *f*) reaction

The results using mono-energetic neutron sources are obviously higher than those using white neutron sources in the MeV region.



**Fig.1** Experimental data from 1.5 to 15 MeV of  $^{232}$ Th(*n*, *f*) reaction using white neutron sources and mono-energetic neutron sources. The results were taken from EXFOR library (after 1970s).

#### **1.4 Differences of** (*n*, *f***) cross-section results from different neutron sources**

**Table I.** The Percentage which the (n, f) cross-section results from monoenergetic neutron sources higher than the results from white neutron sources

Energy region				
Percentage	0.1 - 3.5 MeV	3.5 - 10 MeV	10 - 20 MeV	0.1 - 20 MeV
Nucleus				
<sup>233</sup> U	-1.76%	-1.66%	1.91%	-1.70%
<sup>236</sup> U	-1.29%	0.36%	2.79%	-0.45%
<sup>237</sup> Np	-3.39%	-3.79%	-3.82%	-3.68%
<sup>238</sup> U	-2.84%	0.11%	0.03%	-1.58%
<sup>239</sup> Pu	-1.34%	3.70%	-0.05%	-0.32%
<sup>240</sup> Pu	-0.62%	0.60%	-0.50%	-0.26%
<sup>235</sup> U	2.20%	-1.80%	-1.13%	1.33%
<sup>232</sup> Th	-0.51%	4.64%	-0.61%	2.07%

Data were taken from EXFOR (after 1970s)

The measured  $^{232}$ Th(*n*, *f*) cross-section results using mono-energetic neutron sources are obviously higher than those using white neutron sources especially in the 3.5 - 10 MeV region.

### **1.5 Future Plan**

In order to clarify the existing discrepancies, accurate  ${}^{232}$ Th(*n*, *f*) cross-section measurements in the 3.5 – 10 MeV region are being planned.

### **1.6 Current progress**

Determination of the <sup>232</sup>Th nucleus number by small solid angle method.

(The emitted  $\alpha$  particles with small solid angles would have shorter track in the material, which lead to less energy loss and lower tailing of the peaks.)

### **1.6 Current progress**



**Fig.2** Picture of the <sup>232</sup>Th(OH)<sub>4</sub> sample with tantalum film

- The <sup>232</sup>Th(OH)<sub>4</sub> sample was prepared by Frank Laboratory of Neutron Physics.
- The spectrum from  $^{232}$ Th would be complicated because of the  $\alpha$  particles from the daughter nuclides of  $^{232}$ Th (the small solid angle method could reduce the tailing and overlap of peaks).



Fig.3 Decay chain of <sup>232</sup>Th

### 2.1 The small solid angle device

2.1.1 Compositions of the small angle device

The vacuum chamber;
The vacuum pumping system;
The electronics;
The digital data-acquisition (DAQ) system.



Fig.4 Picture of the device.



### 2.1 The small solid angle device

#### 2.1.2 Compositions of the vacuum chamber

- 1) Three anti-scattering baffles;
- (2) A sample holder;
- 3 A diaphragm;
- 4 A Au-Si surface barrier semiconductor detector;
- 5 Three screw rods;
- 6 Screws and nuts.

(The sample holder, detector, diaphragm and anti-scattering baffles are connected by three screw rods which are mounted at the lid of vacuum chamber, and all of them are centered at a common symmetry axis.)



Fig.5 CAD structure of the vacuum chamber.



Fig.6 Picture of the structure inside the vacuum chamber.

2.2 Geometry

The diameter of diaphragm: 7.00 mm;
The diameter of <sup>232</sup>Th(OH)<sub>4</sub> foil sample : 44.0 mm;
The distance from sample to diaphragm: 29.90 mm.
(the sample-to-diaphragm distance (29.90 mm) = the holder-to-diaphragm distance (30.01 mm) - the height of tantalum (0.11 mm))

(The positions of diaphragm and sample holder were fixed by nuts on the screw rods, and an electronic digital caliper measured the distance.)

### 2.3 The electronics and the DAQ system



Fig.5 Block diagrams of the electronics and the DAQ system

Ps: A data-acquisition software was developed using the LABVIEW language to control the DAQ system, process the measured results and save the waveform data onto the hard disk of our PC.

### 2.4 The experimental process:

1) Foreground measurement for  $\alpha$  events from <sup>232</sup>Th for 1258384s;

2) Background measurement using tantalum film for 601128s.

# 3. Results

### 3.1 Spectra from measurements



#### Characteristics

(1) The background is quiet weak;

(2) The tailing of the peaks from the daughter nuclides would increase the total counts of the  $^{232}$ Th peak. (We used an exponential function to fit the tailing then subtracted from  $^{232}$ Th peak region. The fitting counts accounted for 4.81% of the total counts of the  $^{232}$ Th peak.)

# 3. Results

### 3.1<sup>232</sup>Th nucleus number

The calculated equation of  $^{232}$ Th nucleus number N:

$$N = rac{n_{\mathrm{fore}} - n_{\mathrm{back}} imes rac{t_{\mathrm{fore}}}{t_{\mathrm{back}}} - n_{\mathrm{fit}}}{t_{\mathrm{fore}} imes arepsilon imes \lambda}$$

 $n_{\rm fore}$ : the foreground counts in the <sup>232</sup>Th peak region;

 $n_{\text{back}}$ : the background counts in the <sup>232</sup>Th peak region;

 $n_{\rm fit}$ : the fitting counts in the <sup>232</sup>Th peak region;

 $t_{\rm fore}$ : the durations of foreground measurements;

 $t_{\text{back}}$ : the durations of background measurements;

 $\varepsilon$ : the detection efficiency of the small solid angle device calculated by Monte Carlo simulation (energy loss and back-scattering were taken into account), and the value of  $\varepsilon$  is 0.0024723 in our measurements;  $\lambda$ : the  $\alpha$  decay constant of <sup>232</sup>Th;

The calculated <sup>232</sup>Th nucleus number:  $1.85 \times 10^{19}$ 

# 3. Results

### **3.2 Uncertainties**

(1) The uncertainty of geometry ( $\varepsilon$ ): 0.75%;

(2) The uncertainty of Monte Carlo simulation ( $\varepsilon$ ): 0.33%;

(3) The uncertainty of and unevenness of sample edge ( $\epsilon$ ): 0.30%;

(4) The uncertainty of <sup>232</sup>Th peak region determination ( $n_{fore}$ ): 0.19%;

(5) The uncertainty from the statistical error  $(n_{fore})$ : 0.35%;

(6) The uncertainty of fit  $(n_{\text{fit}})$ : 0.96%;

(7) The uncertainty of the  $\alpha$  decay constant of <sup>232</sup>Th ( $\lambda$ ): 0.40%.

(Other error sources are very small so we ignored them.)

The total uncertainty: 1.41%

### 4. Summary

- The  ${}^{232}$ Th(*n*, *f*) cross-section data from EXFOR library were collected and analyzed, and the measurement results using mono-energetic neutron sources are obviously higher than those using white neutron sources.
- A small solid angle device was designed and installed.
- The <sup>232</sup>Th nucleus number was determined by small solid angle method, and the measured result is 1.85×10<sup>19</sup>(1±1.41%).

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# THANKS